

1 Possible Superconductivity in TiSe_2 Intercalated by 2 Transition Metals

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P1

5 Electric conductivity and magnetoresistance data for TiSe_2 intercalated by Fe^{2+} and Cr^{3+} are
6 interpreted as superconducting state with $T_c \sim 5$ and 4 K, respectively.

7 **KEY WORDS:** superconductivity; intercalation compounds.

8 Intercalation of layered titanium dichalco-
9 genides by silver and transition metals leads to de-
10 crease in the concentration of free charge carriers
11 [1], lattice compression along normal direction to
12 basal plane [2], growth of Ti oxidation state [3].
13 Such a behavior cannot be explained using the rigid
14 band model, usually applied for intercalation com-
15 pounds. These properties were interpreted as a result
16 of appearance of covalent centers Ti–M–Ti, form-
17 ing upon intercalation due to hybridization of $\text{Ti}3d_{z^2}$
18 states of the host lattice and valent states of in-
19 tercalated metal [4]. These centers act as traps for
20 free electrons, while the change of the lattice di-
21 mensions at intercalation was well described as a
22 result of substitution of Ti–M–Ti with Ti–V–Ti (V-
23 vacancy) centers [2]. As lattice deformation and
24 degree of localization of charge carriers at for-
25 mation of the Ti–M–Ti centers are in linear rela-
26 tion, so these centers may be considered as small
27 polarons.

28 As polaron type of charge carriers is supposed
29 to be important for properties of high temperature
30 superconductors (see, for example, [5]), comparison
31 of electrical properties of intercalation compounds
32 with polaron charge carriers at cooling and HTSC
33 cuprates is interesting. For this, we chose Cr_xTiSe_2
34 and Fe_xTiSe_2 systems, which demonstrate relatively

moderate degree of localization, caused by moder- 35
ate polaron shift value [2]. Magnetic moment of in- 36
tercalated impurity has a specific interest; both these 37
systems have antiferromagnetic type of interaction 38
between intercalated atoms, Neel temperature is de- 39
termined by impurity content [6, 7]. 40

Single crystals, grown by gas transport method 41
[3], had the following size: $7 \times 5 \times 0.1 \text{ mm}^3$ for 42
 Fe_xTiSe_2 and $6 \times 5 \times 0.05 \text{ mm}^3$ for Cr_xTiSe_2 . Inter- 43
calant content was estimated using concentration de- 44
pendence of unit cell parameters in *c*-direction, ob- 45
tained earlier for powder samples: $\text{Cr}_{0.33}\text{TiSe}_2$ and 46
 $\text{Fe}_{0.5}\text{TiSe}_2$. Electrical conductivity was measured by 47
standard 4-probe technique in temperature range 48
1.7–300 K and magnetic field $B = 0\text{--}12 \text{ T}$, Center for 49
Magnetic Measurements of Institute of Metal Physics 50
UrD RAS. 51

Relative electrical resistivity for $\text{Fe}_{0.5}\text{TiSe}_2$ is 52
shown in Fig. 1. The anomaly near $\sim 100 \text{ K}$ coin- 53
cides with Neel temperature for the same composi- 54
tion [6] and confirms determined iron content. Rapid 55
drop of resistivity is visible below $\sim 5 \text{ K}$ for $\text{Fe}_{0.5}\text{TiSe}_2$ 56
and $\sim 4 \text{ K}$ for $\text{Cr}_{0.33}\text{TiSe}_2$ (Fig. 2). At the point of 57
the drop the slope of $R/R_{300}(T)$ increases 60 times 58
(Fig. 1). Observed feature looks similar with transi- 59
tion to superconducting state. Magnetoresistivity 60
 $\rho(B)$ for both systems is shown in Fig. 3. This value 61
increases as B increase and demonstrates a kink at 62
some critical magnetic field 0.46 T for $\text{Fe}_{0.5}\text{TiSe}_2$ and 63
0.40 T for $\text{Cr}_{0.33}\text{TiSe}_2$ which we attribute with sup- 64
pression of superconductivity. At higher magnetic 65
field up to 12 T for $\text{Fe}_{0.5}\text{TiSe}_2$ $\rho(B)$ remains linear 66
(Fig. 4). 67

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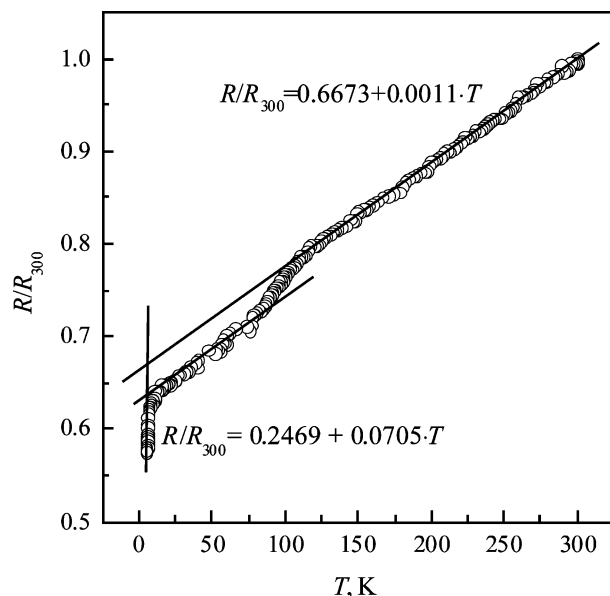
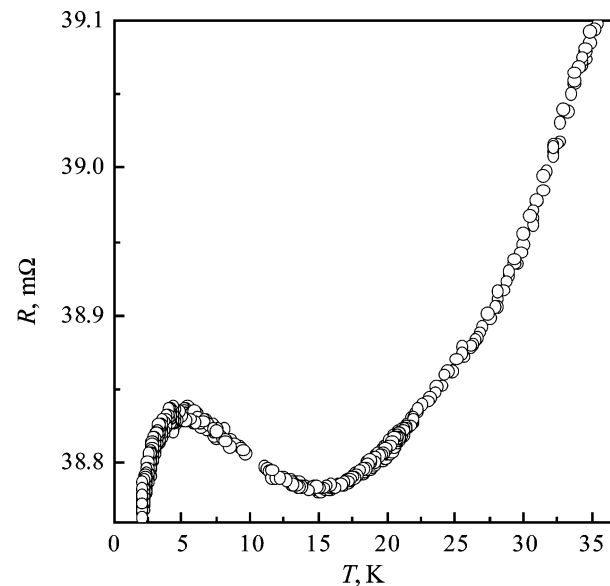


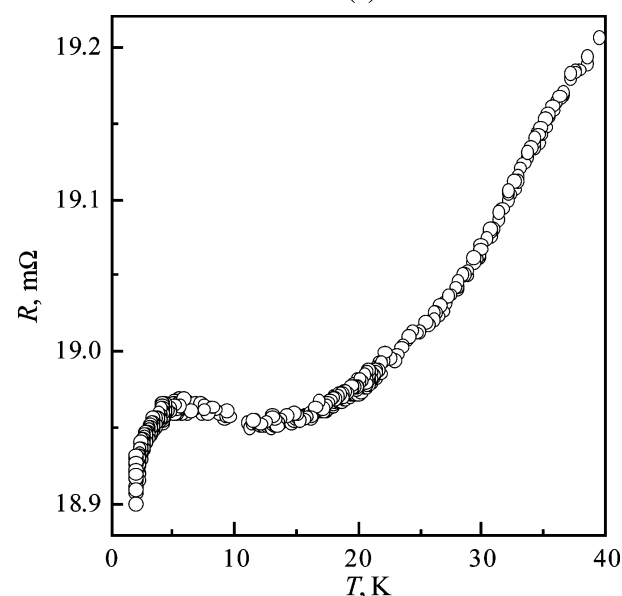
Fig. 1. Relative electrical resistivity for $\text{Fe}_{0.5}\text{TiSe}_2$.

68 Unfortunately, we did not have the possibility
 69 to decrease the temperature below 1.7 K and thus,
 70 were not able to observe the transition in whole. The
 71 observed effect is small, moreover, linear extrapolation
 72 of $R(T)$ in “superconducting” state (Fig. 1)
 73 gives a positive remnant resistivity at $T = 0$. It
 74 may be explained as inhomogeneous state of the
 75 material, when superconducting fractions are alter-
 76 nated by normal metal. Another explanation is the
 77 small concentration of Cuper pairs. Temperature de-
 78 pendences of electrical resistivity do not show any
 79 anomaly, which may be attributed with phase separa-
 80 tion. Moreover, the phase transition of first order
 81 has kinetic barrier at low temperature because of low
 82 diffuse mobility of the components. So the first expla-
 83 nation seems less probable.

84 Small Cuper pair concentration may have the
 85 following reason. As it was predicted from band
 86 calculations, the hybrid state band $\text{Ti}3d_{z^2}/\text{M}3d$ for
 87 $\text{M} = \text{Cr}, \text{Fe}$ should have spin splitting [8]. For
 88 Cr_xTiSe_2 such a splitting was experimentally ob-
 89 served [3]. Evidently, when the splitting becomes sig-
 90 nificant it is possible to observe the gap between
 91 spin sub-bands. Observed activation type of electri-
 92 cal conductivity for $\text{Cr}_{0.33}\text{TiSe}_2$ [3] corresponds to
 93 this situation. For Cr^{3+} $x = 0.33$ concentration cor-
 94 responds to transfer of one electron per unit cell,
 95 or by other words, exactly half filling of the hy-
 96 brid state band. For Fe^{2+} -intercalated compound,



(a)

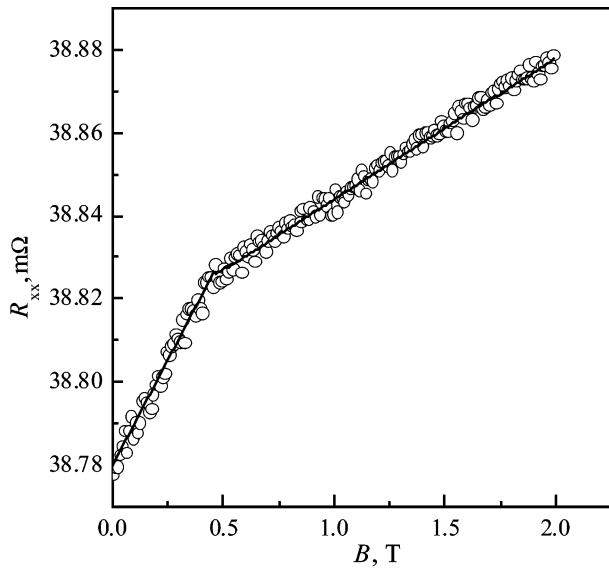


(b)

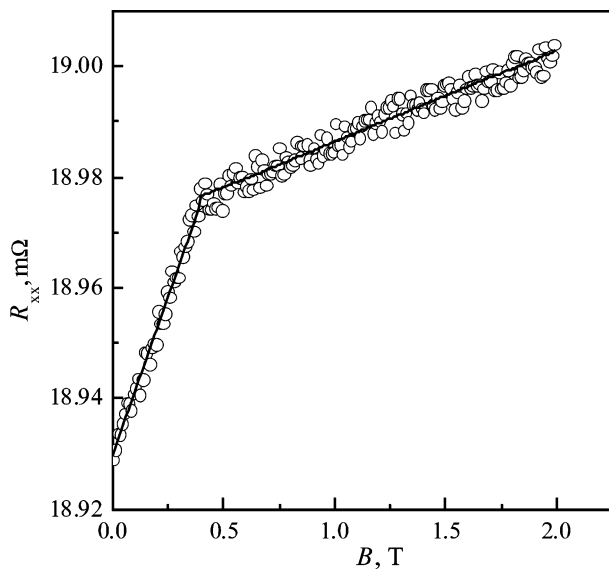
Fig. 2. Electrical resistivity for $\text{Fe}_{0.5}\text{TiSe}_2$ (a) and $\text{Cr}_{0.33}\text{TiSe}_2$ (b). Drop of resistivity below ~ 5 K for $\text{Fe}_{0.5}\text{TiSe}_2$ and ~ 4 K for $\text{Cr}_{0.33}\text{TiSe}_2$ we connect with superconducting state.

the same situation will be obtained at $x=0.5$. The
 investigated single crystals had exactly these com-
 positions what makes superconductivity impossible.
 In real crystals small amount of defects is possi-
 ble, enough to form small quantity of Cuper pairs.
 These defects must have some excess of metal in
 comparison with stoichiometric composition (when

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(a)



(b)

Fig. 3. Electrical resistivity for $\text{Fe}_{0.5}\text{TiSe}_2$ (a) and $\text{Cr}_{0.33}\text{TiSe}_2$ (b) as a function of external magnetic field. The kink at $B = 0.46$ T for $\text{Fe}_{0.5}\text{TiSe}_2$ and $B = 0.40$ T for $\text{Cr}_{0.33}\text{TiSe}_2$ we connect with suppress of superconductivity.

104 the Fermi level is located into low spin sub-band, all
 105 electrons have the same spin and Cuper pairs can
 106 not be formed). This supposition may explain an ab-
 107 sence of Meissner effect for powder $\text{Cr}_{0.5}\text{TiSe}_2$ [7],
 108 where the samples with precise stoichiometric com-
 109 position were studied. With an aim to check this
 110 supposition, the measurement of temperature de-

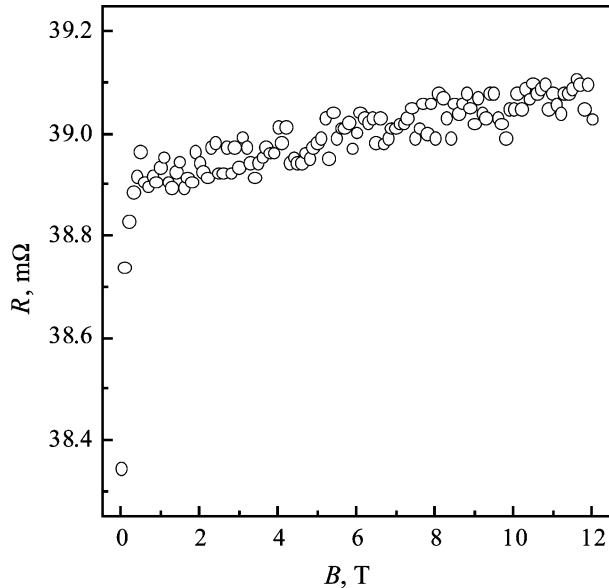


Fig. 4. Electrical resistivity of $\text{Fe}_{0.5}\text{TiSe}_2$ as a function of external magnetic field. The linear behavior remains up to 12 T.

pendence of electric conductivity for powder sample 111
 with $\text{Fe}_{0.5}\text{TiSe}_2$ precise stoichiometric composition 112
 was done. The method of sample preparation and 113
 measurement was the same as for Cr_xTiSe_2 [3]. The 114
 data, shown in Fig. 5, demonstrate activation type 115

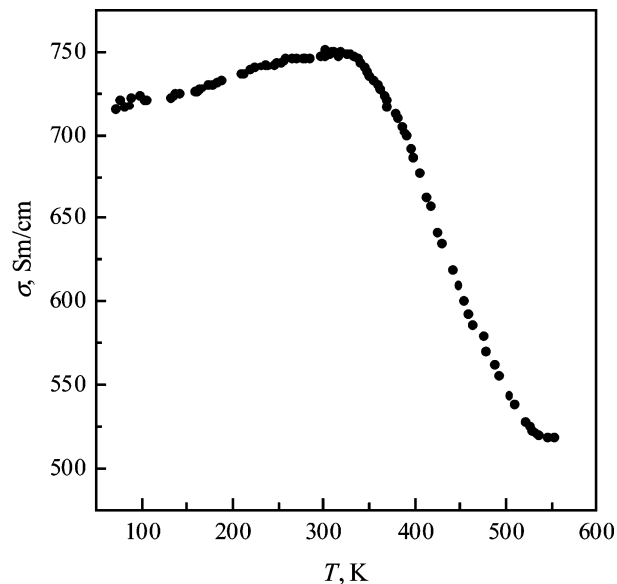


Fig. 5. Temperature dependence of electrical resistivity for powder sample with $\text{Fe}_{0.5}\text{TiSe}_2$ precise composition. Activation type of low temperature part of this function shows the closeness of Fermi level to energy gap.

116 of electric conductivity, in agreement with above
 117 supposition.
 118 The relation between over-stoichiometric de-
 119 fects and superconducting state remains unclear and
 120 requests following investigation.

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